

Inertia and Chiral Edge Modes of a Skyrmion Magnetic Bubble

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(Received 21 August 2012; published 20 November 2012)

The dynamics of a vortex in a thin-film ferromagnet resembles the motion of a charged massless particle in a uniform magnetic field. Similar dynamics is expected for other magnetic textures with a nonzero Skyrmion number. However, recent numerical simulations reveal that Skyrmion magnetic bubbles show significant deviations from this model. We show that a Skyrmion bubble possesses inertia and derive its mass from the standard theory of a thin-film ferromagnet. In addition to center-of-mass motion, other low energy modes are waves on the edge of the bubble traveling with different speeds in opposite directions.

DOI: [10.1103/PhysRevLett.109.217201](https://doi.org/10.1103/PhysRevLett.109.217201)

PACS numbers: 75.70.Kw, 75.78.-n

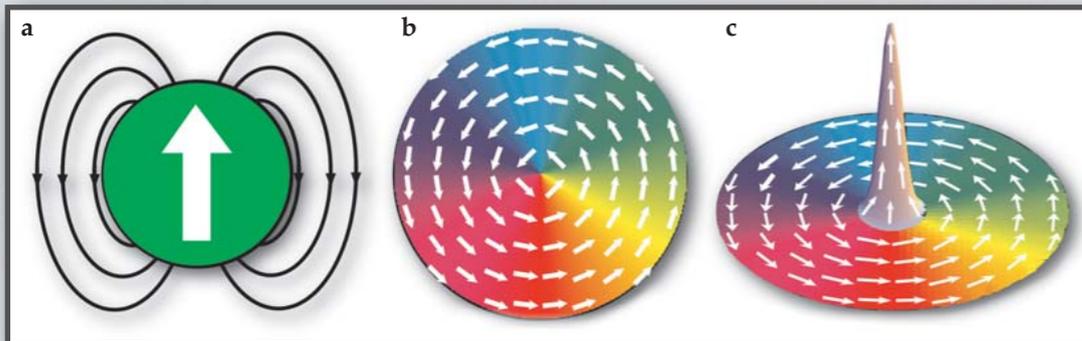
TOPOLOGICAL STATES IN FM NANOMAGNETS

Nanomagnets

Disk of 15 nm thickness and $D=100$ nm diameter

Competition between

- Exchange interaction (wants to align all spins, density decreases with D)
- Magnetic stray field (wants to create vortex-like structure, increases with D)



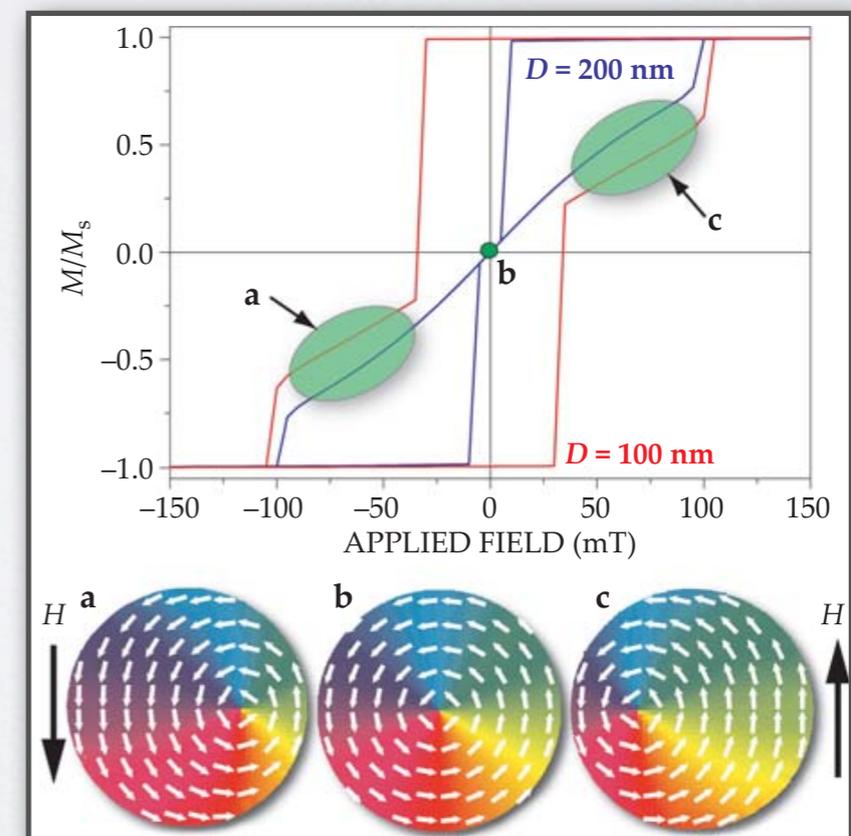
Magnetic field

Due to Zeeman interaction spins want to align with B

For intermediate magnetic fields (up to ~ 50 - 300 mT) vortex structure stays in tact

Allows us to move vortex around in a controlled way

C. L. Chien *et al.*, *Physics Today* **60**(6), 40 (2007)



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Interest: classical magnetic memories (MRAM)

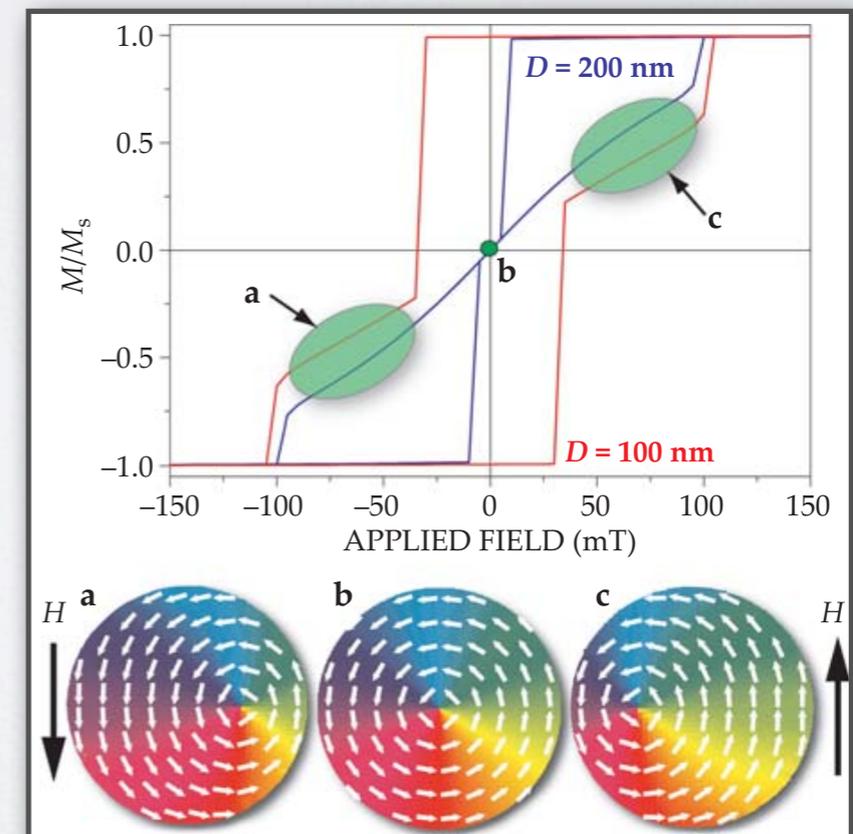
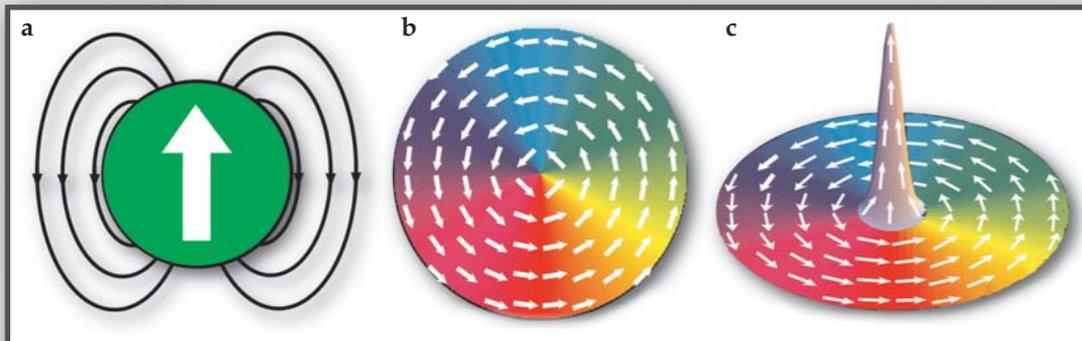
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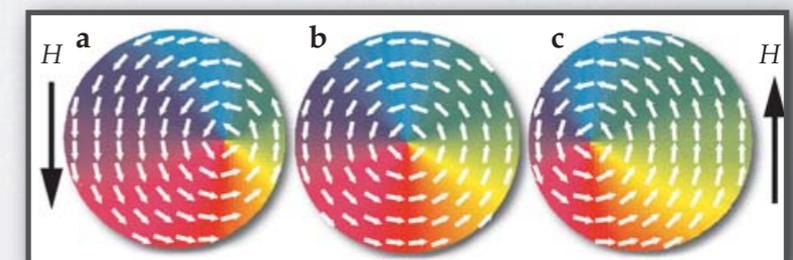
A. A. Thiele, Phys. Rev. Lett. **30**, 230 (1973)

How can we describe the motion of such **rigid** magnetization structures?

Start from the Landau-Lifshitz equation:

$$\dot{\mathbf{m}} = \underbrace{\gamma \mathbf{B} \times \mathbf{m}}_{\text{Precession}} + \underbrace{\alpha \mathbf{m} \times \dot{\mathbf{m}}}_{\text{Damping}}$$

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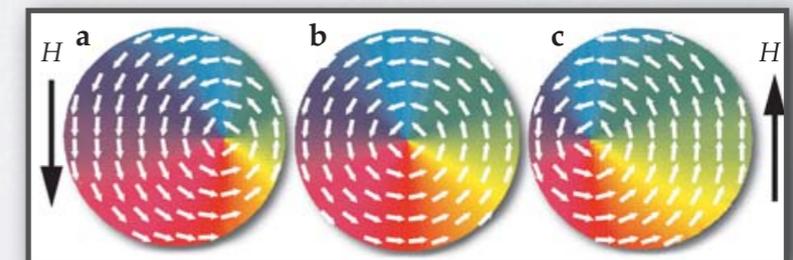
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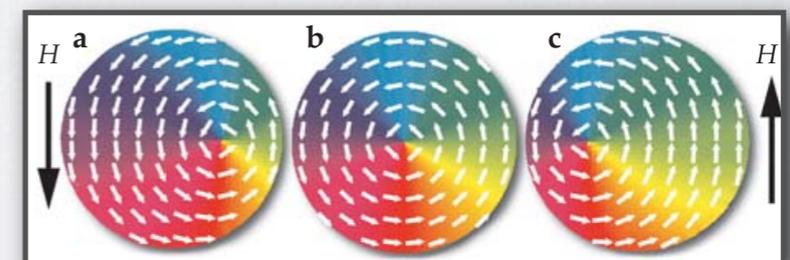
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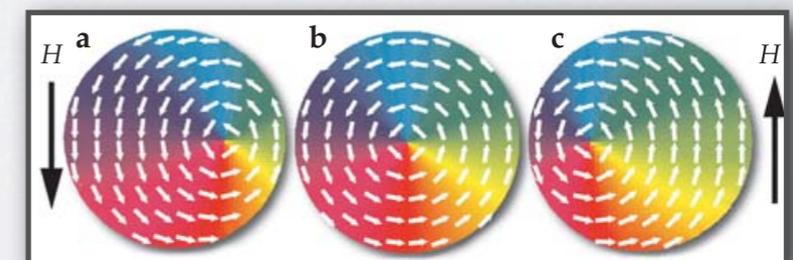
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$$\mathbf{G} \times \dot{\mathbf{R}} + \mathbf{F} - D\dot{\mathbf{R}} = 0$$



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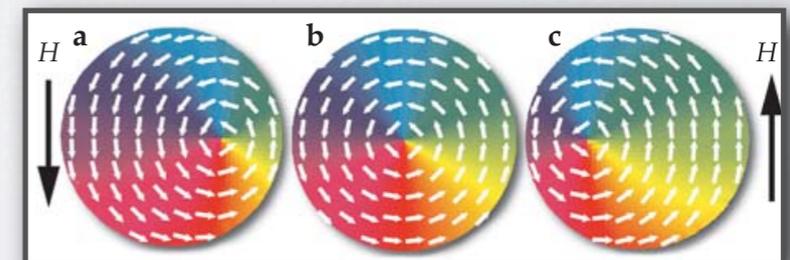
$$\mathbf{G} \times \dot{\mathbf{R}} + \mathbf{F} - D \dot{\mathbf{R}} = 0$$

The parameters \mathbf{G} and D depend on the magnetization profile, for vortex

$$\mathbf{G} = (0, 0, \mathcal{G}) \quad \text{where} \quad \mathcal{G} \propto (1/4\pi) \int d^2 \mathbf{r} \mathbf{m} \cdot (\partial_x \mathbf{m} \times \partial_y \mathbf{m})$$



Topological invariant, Skyrmion number $q=1/2$



THIELE'S EQUATION (2)

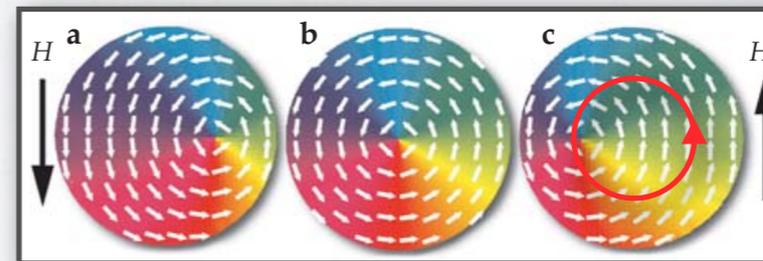
$$\mathbf{G} \times \dot{\mathbf{R}} + \mathbf{F} - D\dot{\mathbf{R}} = 0$$

Corresponds to the EOM of a massless particle in a magnetic field \mathbf{G}/e and force $\mathbf{F} = -\nabla U$

$$\begin{array}{cccc}
 -M\ddot{\mathbf{R}} & + & e\mathbf{B} \times \dot{\mathbf{R}} & + & \mathbf{F} & - & D\dot{\mathbf{R}} & = & 0 \\
 \underbrace{\phantom{-M\ddot{\mathbf{R}}}} & & \underbrace{\phantom{e\mathbf{B} \times \dot{\mathbf{R}}}} & & \underbrace{\phantom{\mathbf{F}}} & & \underbrace{\phantom{D\dot{\mathbf{R}}}} & & \\
 \downarrow & & \downarrow & & \downarrow & & \downarrow & & \\
 \text{Mass } (=0) & & \text{Lorentz} & & \text{External} & & \text{Damping} & & \\
 & & \text{Force} & & \text{Force} & & (=0) & &
 \end{array}$$

Example: Motion of a magnetic vortex in a parabolic potential well (ignore damping)

$$\begin{array}{c}
 U(X, Y) = \mathcal{K}(X^2 + Y^2)/2 \\
 \underbrace{\phantom{U(X, Y) = \mathcal{K}(X^2 + Y^2)/2}} \\
 \downarrow \\
 \text{Due to 'magnetostatic} \\
 \text{interaction between} \\
 \text{vortex and edge of disc'}
 \end{array}$$



Leads to circular motion with frequency: $\omega = \mathcal{K}/\mathcal{G}$

MOTION OF SKYRMIONS (MAGNETIC BUBBLES)

Skyrmions are similar to vortices, but with skyrmion number $q=1$



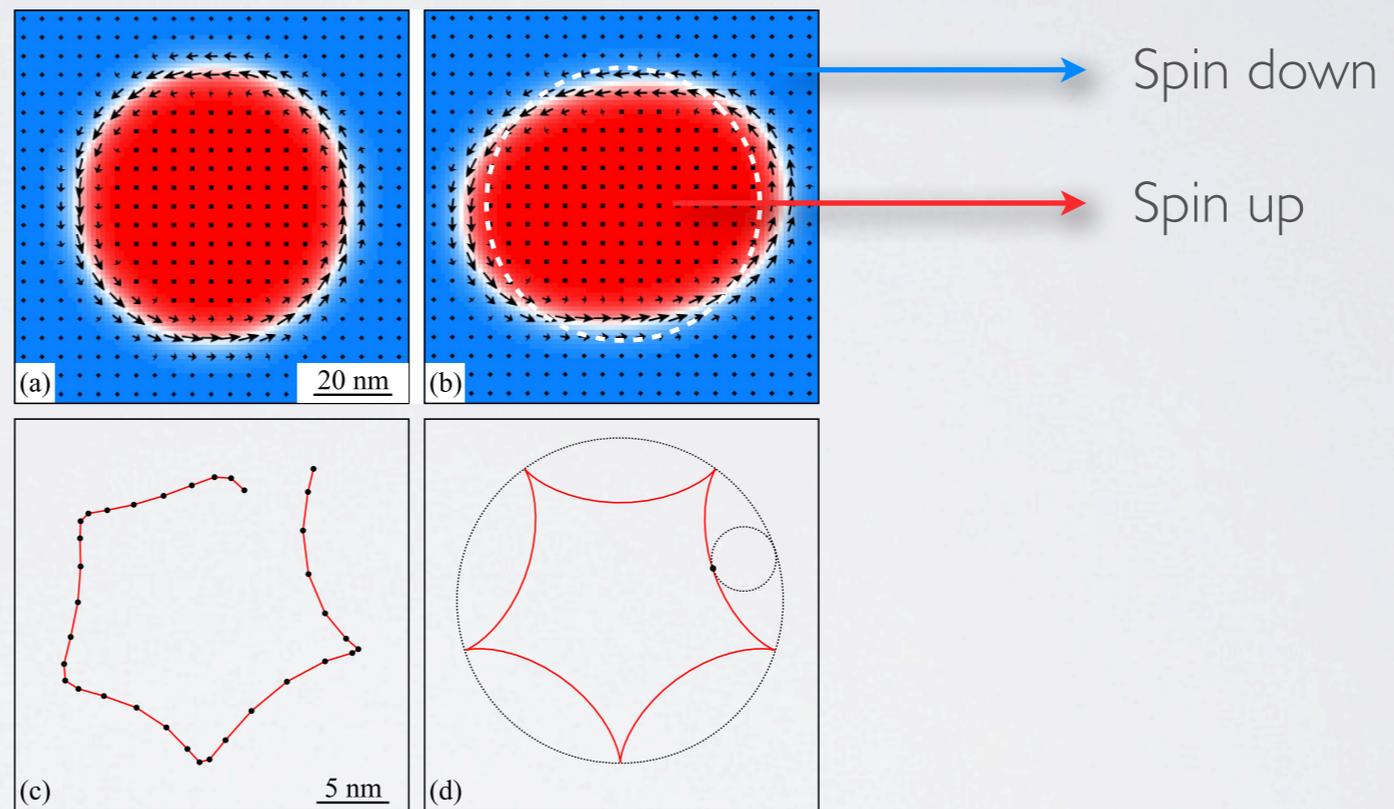
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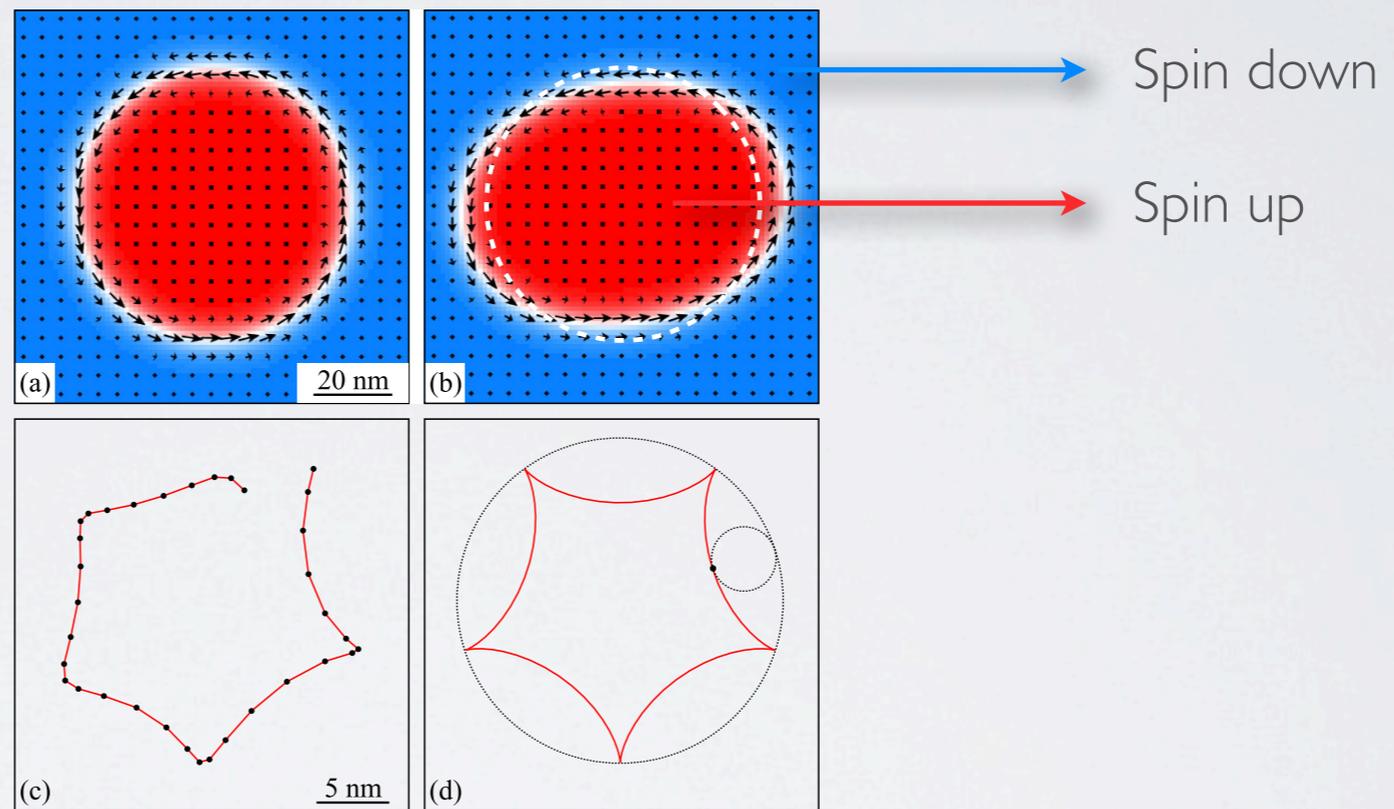
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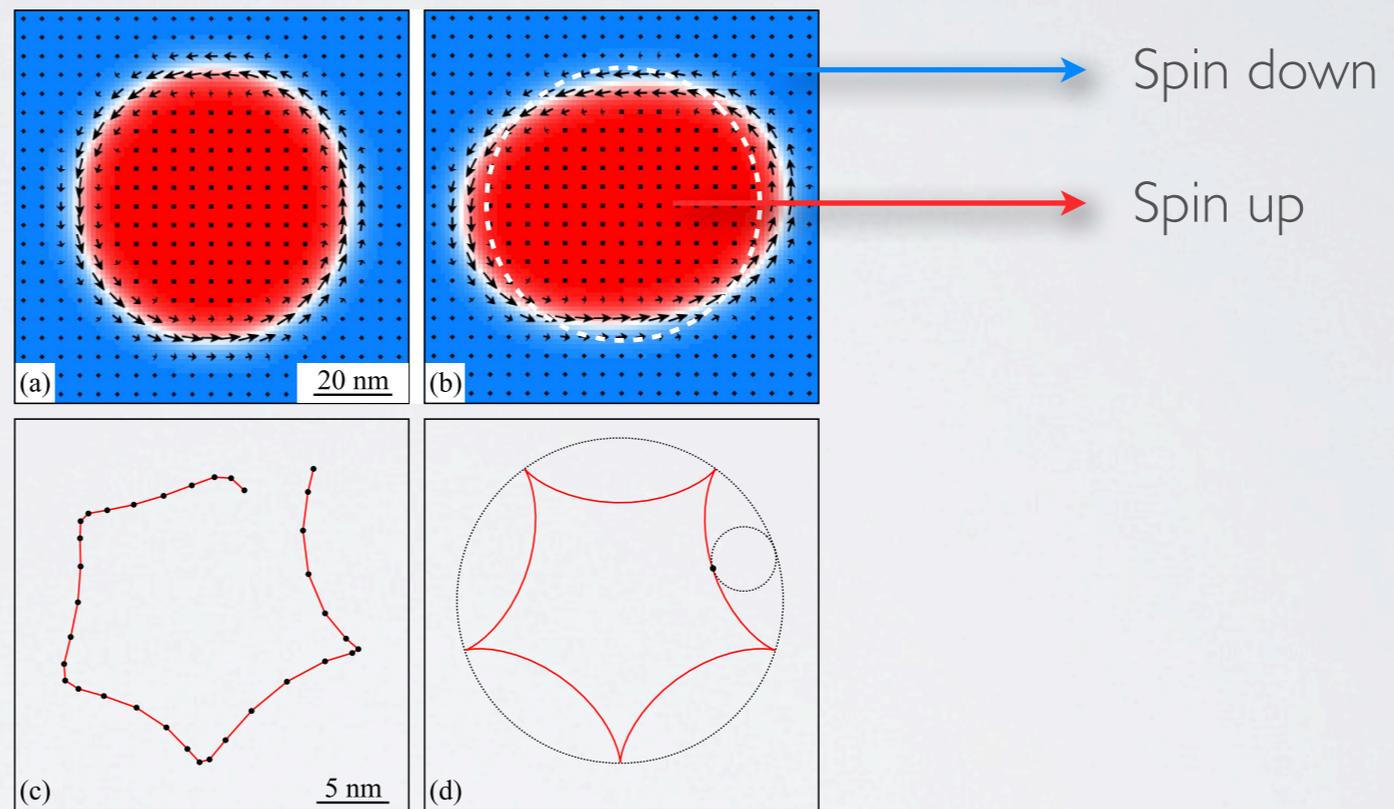
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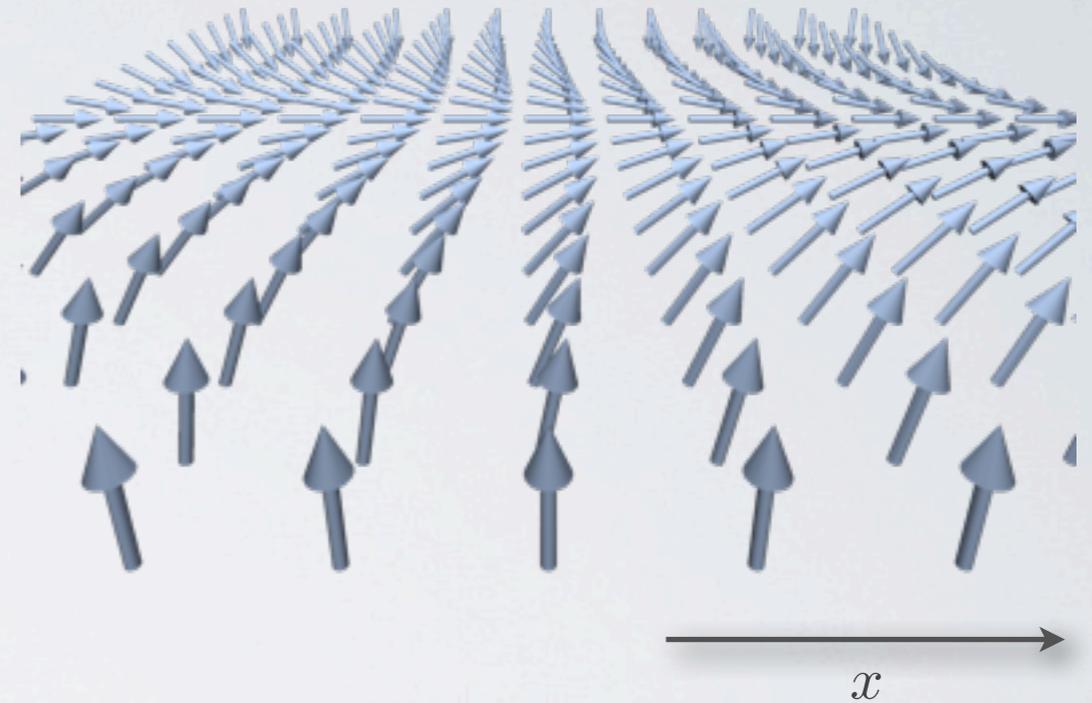
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What is the physical origin of this mass?

ORIGIN OF MASS TERM: STRAIGHT DOMAIN WALL (I)

Certain small **deformations** of the magnetic texture turn out to generate the mass-term



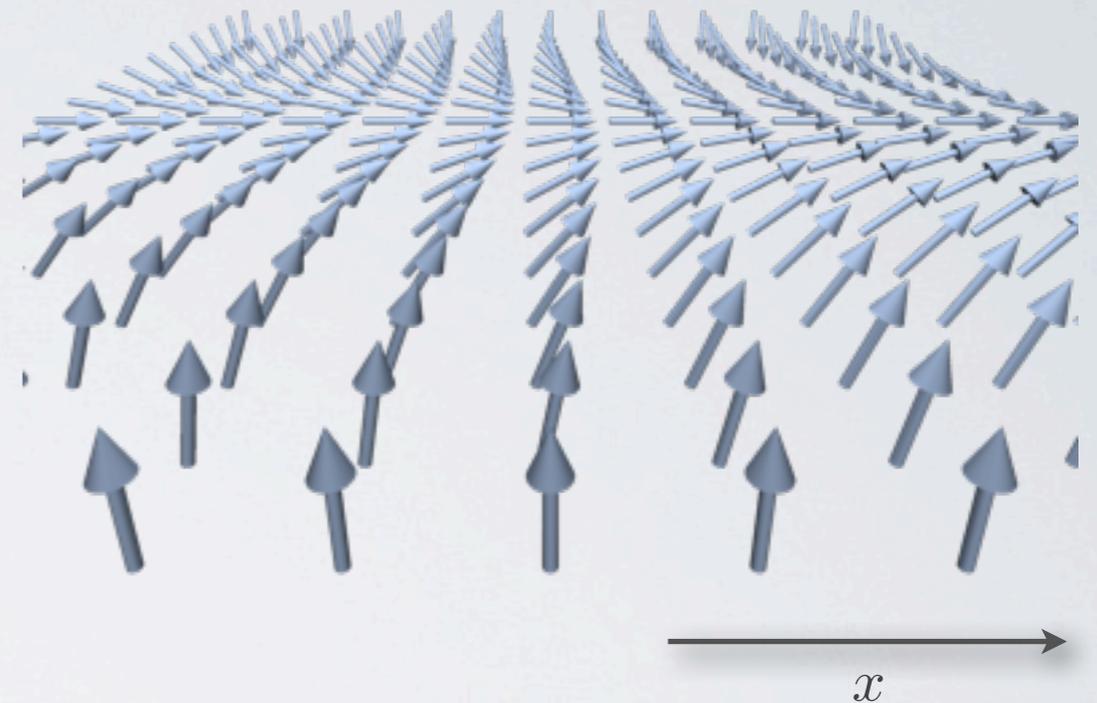
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In the Landau-Lifshitz formalism:

$$\mathbf{m} = (\sin \theta \cos \psi, \sin \theta \sin \psi, \cos \theta)$$

$$L[\theta, \psi] = \int d^2\mathbf{r} \left\{ (1 - \cos \theta)\dot{\psi} - U[\theta, \psi] \right\}$$



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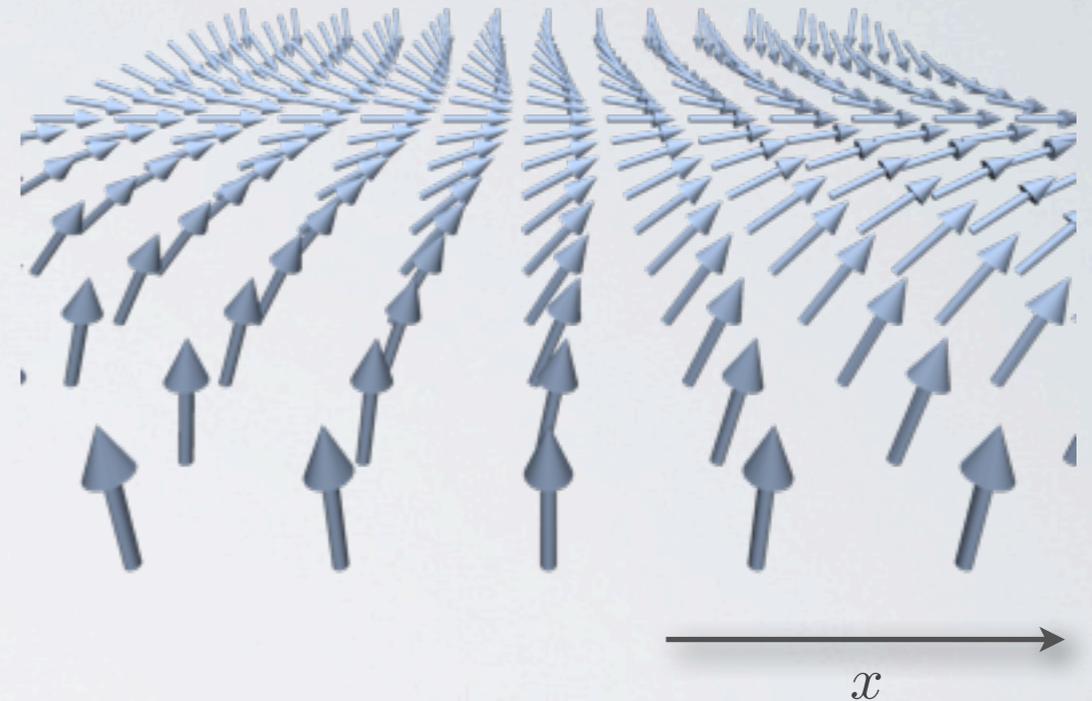
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For this specific case, $\theta = \theta(y, Y)$. $Y(x, t)$ describes the line where $m_z=0$. New variables

$$\psi(x, t) \text{ and } Y(x, t)$$



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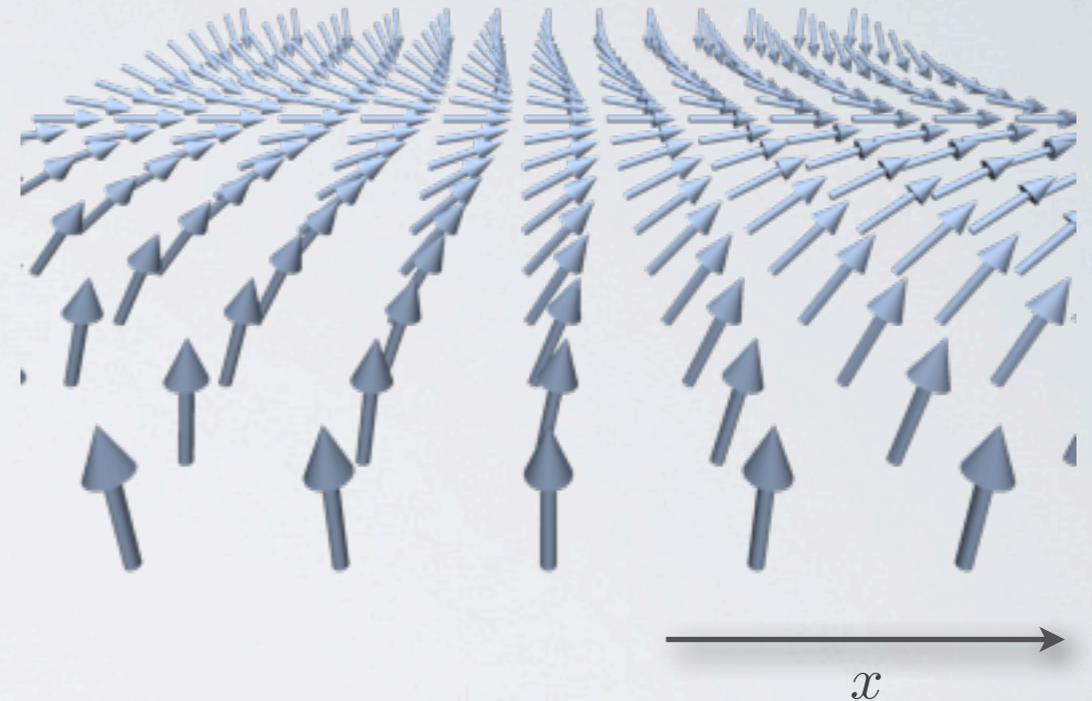
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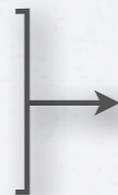
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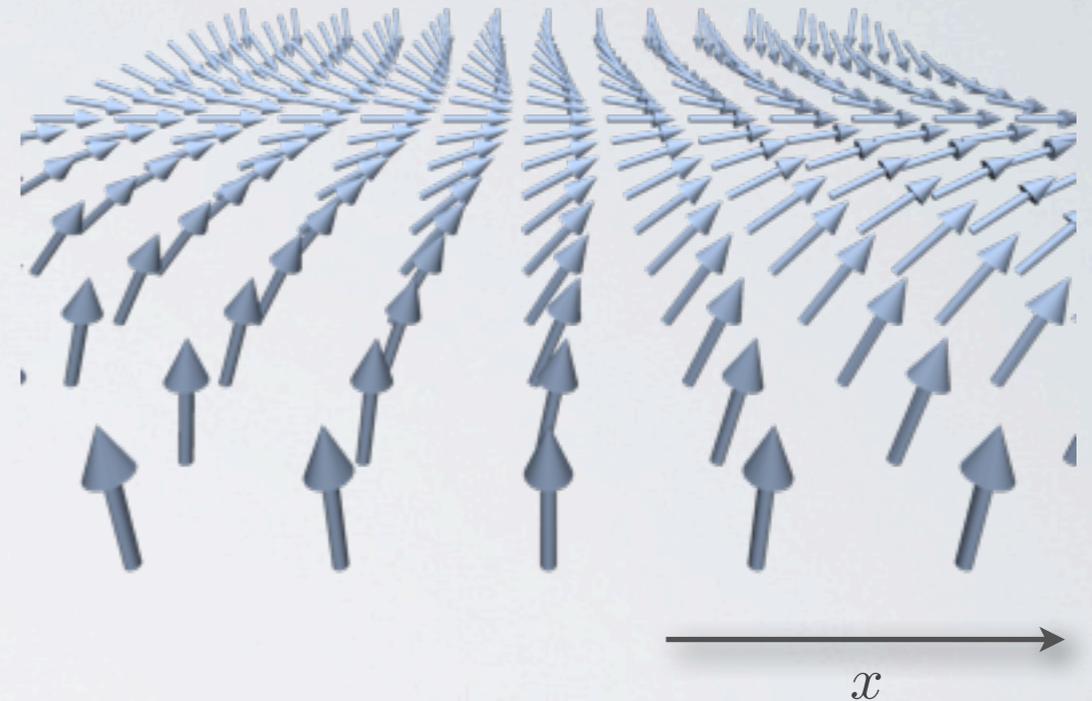
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Gives the two equations of motion:

$$\begin{aligned} -g \dot{\psi} - \delta U / \delta Y &= 0 \\ g \dot{Y} - \delta U / \delta \psi &= 0 \end{aligned}$$



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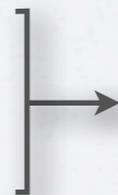
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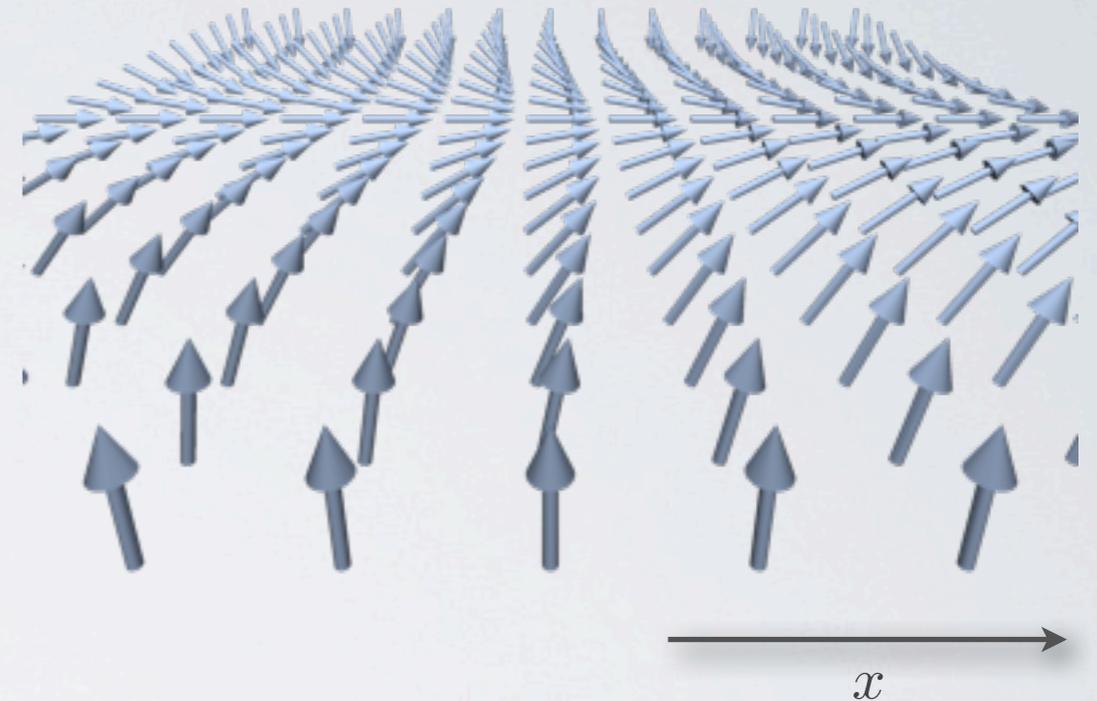


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In equilibrium: in plane magnetization is **aligned with the wall**

$$\psi = Y' \equiv \partial Y / \partial x$$



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Magnetization can be **misaligned** with the wall. Costs an energy:

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Potential energy. Two contributions:

- Local
- Long-range

Local interaction given by

$$U_l[Y] = \int \sigma \sqrt{X^2 + Y^2} \approx U_l[0] + \int_0^l dx \sigma Y'^2 / 2$$

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Waves traveling left and right with velocity

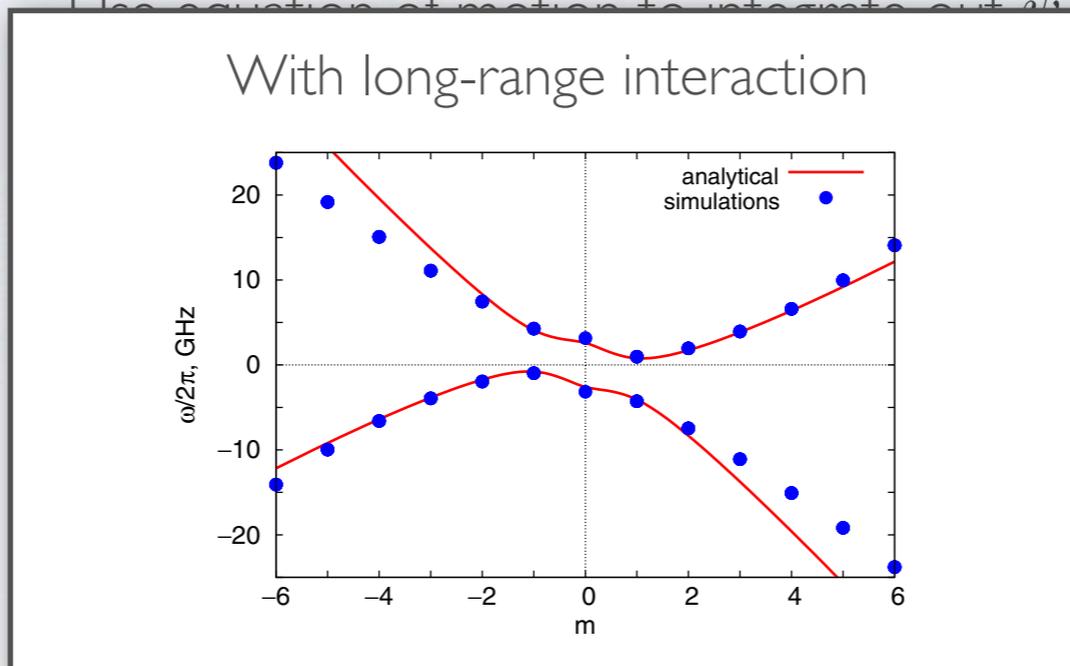
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ORIGIN OF MASS TERM: MAGNETIC BUBBLE

Works pretty much the same as for domain wall.

Parametrization domain wall and magnetization:

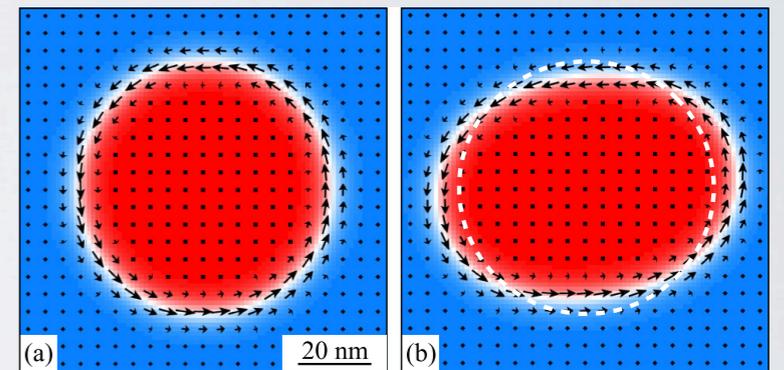
$$r(\phi) = \bar{r} + \sum_m r_m e^{im\phi} \quad \psi(\phi) = \phi \pm \pi/2 + \sum_m \psi_m e^{im\phi}$$

This gives the **Lagrangian**

$$L[r, \psi] = 2\pi\bar{r} \sum_m \left(g\dot{r}_m \psi_m^* - \frac{\mathcal{K}|\psi_m + imr_m/\bar{r}|^2}{2} \right) - U[r]$$

After removing magnetization angle

$$L[r] = \sum_m \left(\pi\bar{r}\rho|\dot{r}_m|^2 - 4\pi m i g r_m^* \dot{r}_m \right) - U[r]$$



Assuming only the lowest mode $r_1 = (X - iY)/2$ is excited

$$L[X, Y] = \mathcal{M}(\dot{X}^2 + \dot{Y}^2)/2 + \mathcal{G}\dot{X}Y - \mathcal{K}(X^2 + Y^2)/2$$

CONCLUSIONS

Thiele's equation predicts that rigid magnetic textures behave as massless particles moving in a magnetic field determined by the texture

This approach works well for magnetic bubbles, but fails for skyrmions

The authors show that the correct motion is predicted when a mass term is added to Thiele's equation

The authors find the origin of this mass term

THANK YOU FOR YOUR ATTENTION